RICH METHANOL COMBUSTION IN SMALL-SCALE COUNTER-FLOW BURNERS TO PRODUCE HYDROGEN-RICH SYNGAS

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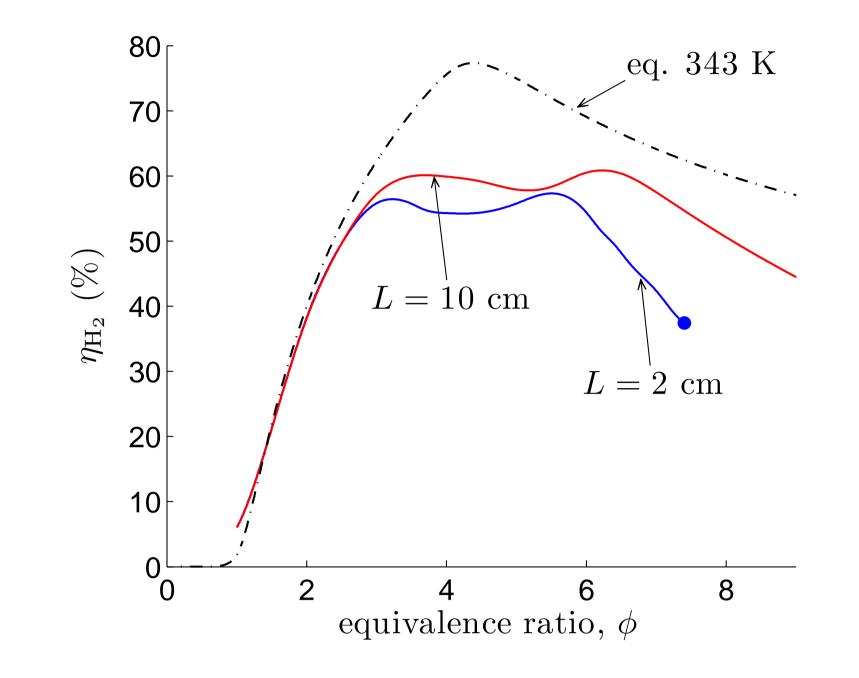
Introduction

In an effort to investigate the suitability of the concept of portable hydrogen production, rich methanol combustion (or partial oxidation) in a small-scale counter-flow burner consisting of an array of multiple narrow channels (see Fig. 1) is examined. The problem is described within the framework of a one-dimensional model (narrowchannel approximation [2]). Methanol oxidation is modeled with detailed chemistry and transport using the recently revised 247-step San Diego mechanism that involves N = 50 reactive species [1].

Results

Two different lengths of the channels L = 2 cm and L = 10 cm are investigated. In Fig. 2 we depict the temperature and hydrogen molar fraction profiles along channel *i* for different equivalence ratios at a fixed inflow velocity. In Fig. 3 we plot them for different inflow velocities at a fixed equivalence ratio. We also plot in the figures (dot-dashed lines) the corresponding values for chemical equilibrium conditions for $\phi = 2, 3$ and 4.

divided by the H_2 moles produced if all the atomic hydrogen bound in the fuel was converted to H_2 . It is calculated as $\eta_{\rm H_2} = \dot{N}_{\rm H_2,out}/2\dot{N}_{\rm CH_3OH,u}$, with \dot{N}_j the molar flow rate of the species j. Maximum efficiency is found for equivalence ratios in the vicinity of 3 and inlet velocity of 200 cm/s.



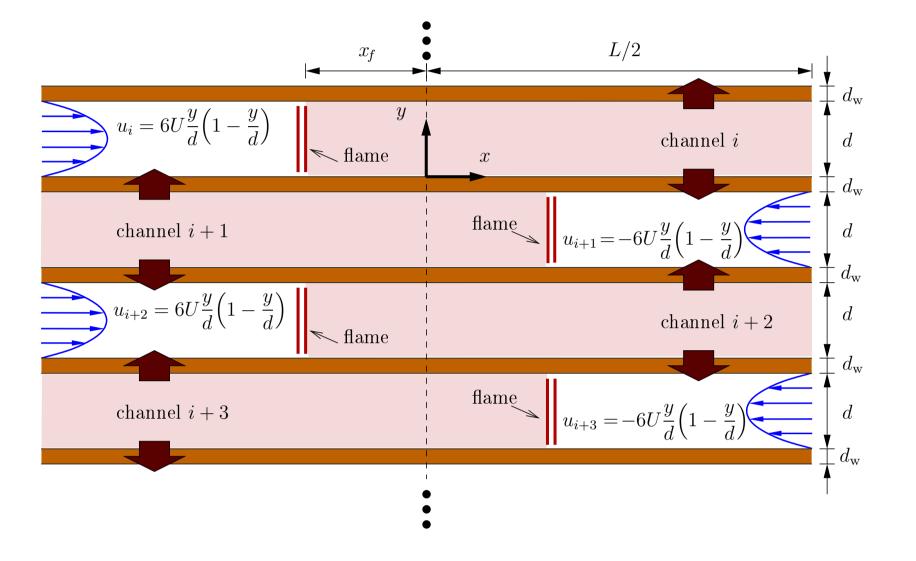


Fig. 1: Sketch of an array of multiple counter-flow narrow channels. L is the length of the channels, d is the distance between the internal surfaces of the walls, $d_{
m W}$ is the thickness of the solid conducting walls, and x_f is half the distance between the opposed flames.

One-dimensional model

Assuming equal gas inflow velocity U_u in all channels (thus flames stabilize at symmetric positions $\pm x_f$ due to heat recirculation [3]), periodicity allows to reduce the analysis to a single channel i with the following governing gas-phase equations

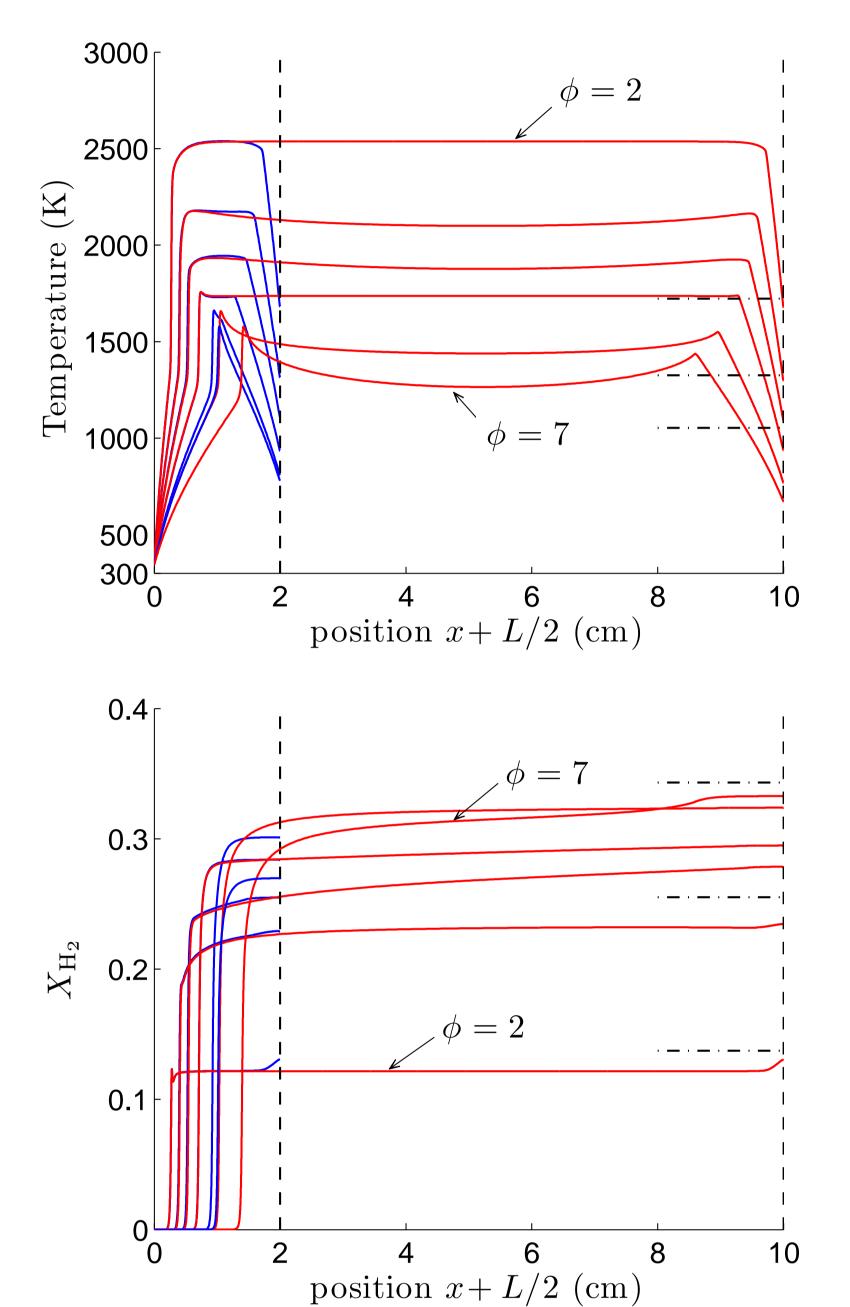
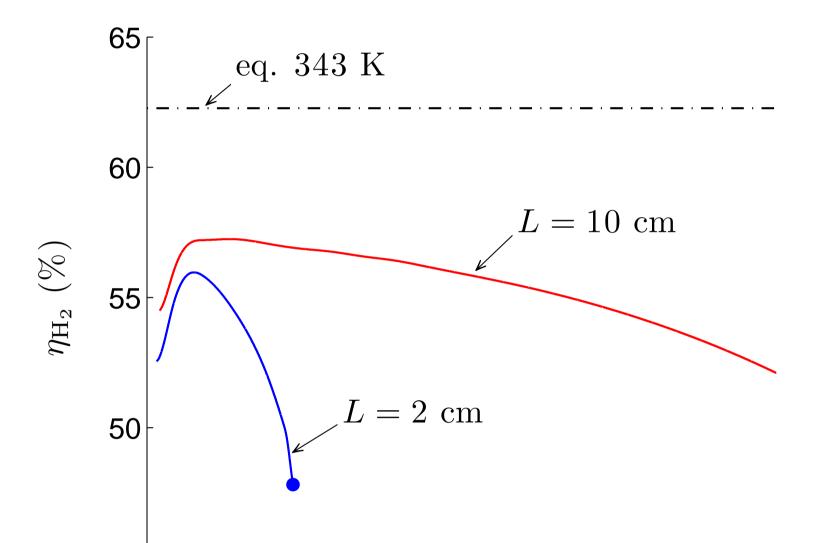


Fig. 4: Methanol-to-hydrogen conversion efficiency variation with the equivalence ratio. Calculated for $U_u = 200 \text{ cm/s}$, $T_u = 343 \text{ K}$, p = 1atm, and $b = 10^6 \text{ W/m}^3 \text{ K}$. The dashed curve corresponds to equilibrium conditions. The solid circle indicates the blow-off limit (beginning of the unstable branch).



$$\rho U = \rho_u U_u = \text{const.},$$

(1)

(3)

(4)

(5)

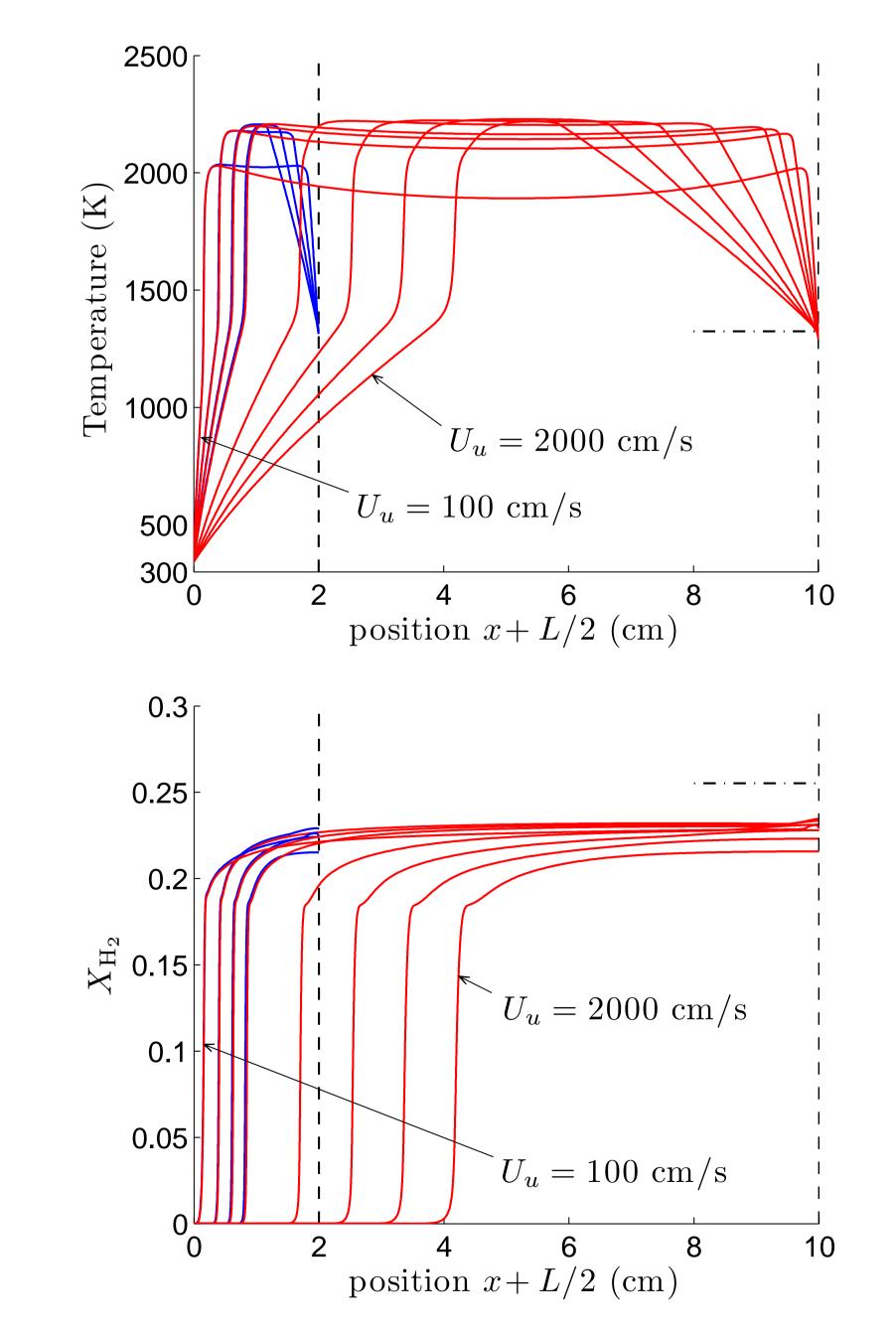
(6)

$$\rho c_p U \frac{\mathrm{d}T}{\mathrm{d}x} = \frac{\mathrm{d}}{\mathrm{d}x} \left(\lambda_g \frac{\mathrm{d}T}{\mathrm{d}x} \right) - \rho \sum_{k=1}^N c_{p,k} Y_k V_k \frac{\mathrm{d}T}{\mathrm{d}x} - \sum_{k=1}^N \dot{\omega}_k h_k W_k - \frac{2\lambda_w}{\mathrm{d}d_w} [T(x) - T(-x)], \quad (2)$$
$$\rho U \frac{\mathrm{d}Y_k}{\mathrm{d}x} = -\frac{\mathrm{d}}{\mathrm{d}x} \left(\rho Y_k V_k \right) + \dot{\omega}_k W_k, \qquad k = 1, 2, \dots, N,$$

 $\rho = \frac{pW}{\mathcal{R}T},$ and with the flux boundary conditions [4] at the left and right ends given by

$$x = -L/2: \begin{cases} (\rho Uh)_u = \rho Uh - \lambda_g \frac{\mathrm{d}T}{\mathrm{d}x} + \sum_{k=1}^N \rho Y_k V_k h_k, \\ (\rho UY_k)_u = \rho UY_k + \rho Y_k V_k, \end{cases}$$

Fig. 2: Profiles of temperature and hydrogen molar fraction along channel *i* for L = 2 cm and L = 10 cm. Calculated for $\phi = \{2, 3, 4, 5, 6, 7\}$, $U_u = 200 \text{ cm/s}, T_u = 343 \text{ K}, p = 1 \text{ atm, and } b = 10^6 \text{ W/m}^3 \text{ K}.$



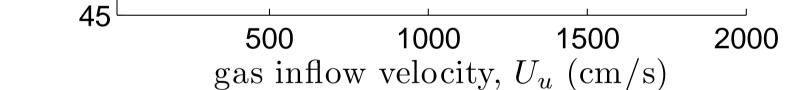


Fig. 5: Methanol-to-hydrogen conversion efficiency variation with the inlet velocity. Calculated for $\phi=3$, $T_u=343$ K, p=1 atm, and $b = 10^6 \text{ W/m}^3 \text{ K}.$

Future work

- Study the effect of external heat losses.
- Study the effect of water addition (autothermal reforming)
- Investigate the effect of flame curvature (2D effects)

References

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$$x = L/2$$
: $\frac{\mathrm{d}T}{\mathrm{d}x} = 0, \ \frac{\mathrm{d}Y_k}{\mathrm{d}x} = 0.$

This one-dimensional formulation is identical to that describing the propagation of a planar flame, with U evaluated at the unburnt conditions being the flame propagation velocity. Here, however, $U(x) = \{\int_0^d u(x, y) \, dy\}/d$ corresponds to the average flow velocity within the prescribed Poiseuille flow. The parameter $b = 2\lambda_w/(d d_w)$ appears in the energy equation as a volumetric heat-recirculation term and measures the amount of heat transferred across the dividing solid walls. For a micro-burner with d = 4mm and $d_{\rm w} = 1$ mm the order of magnitude of b lies typically within 10^5 W/(m^3 K) for materials as quartz to 10^7 $W/(m^{3}K)$ for silicon carbide [5].

Fig. 3: Profiles of temperature and hydrogen molar fraction along channel i for L = 2 cm (dashed curves) and L = 10 cm (solid curves). Calculated for $U_u = \{100, 200, 300, 400, 800, 1200, 1600, 2000\}$ cm/s, $\phi = 3$, $T_u = 343$ K, p = 1 atm, and $b = 10^{6}$ W/m³ K.

Figs. 4 and 5 show the conversion efficiency of methanol to hydrogen, defined, as in [5], as the ratio of H_2 moles produced per unit time at the exit of the reactor

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